which then undergoes an insertion reaction to give the observed $\eta^{3}$-crotyl species. In agreement with this suggestion reaction


$3, \mathrm{~L}=\mathrm{P}(\mathrm{OMe})_{3}$



6. $\mathrm{L}=\mathrm{P}(\mathrm{OMe})_{1}$
with $\mathrm{NaBD}_{4}$ proceeds regioselectively to give 6 with deuterium incorporated in the illustrated position adjacent to the methyl group. A related $\beta$-allylic elimination reaction forming a $\eta^{3}$-crotyl complex has been observed ${ }^{7}$ on heating the $E$ isomer of $\operatorname{Ir}\left[\mathrm{C}(\mathrm{Me})=\mathrm{CHMe}^{2} \mathrm{COL}_{2}\right.$; however, in contrast to the molybdenum system the $Z$ isomer undergoes a cis- $\beta$-vinyl H elimination to form the corresponding hydride and but-2yne.

Acknowledgment. We thank the S.R.C. for support.

## References and Notes

(1) Reaction (room temperature, $\mathrm{CH}_{2} \mathrm{Cl}_{2}$ ) of $\left[\mathrm{Mo}(\mathrm{CO})_{3}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)\right]_{2}$ or [Mo-$\left.(\mathrm{CO})_{3}\left(\eta^{5}-\mathrm{C}_{9} \mathrm{H}_{7}\right)\right]_{2}$ with $\mathrm{AgBF}_{4}$ in the presence of an excess of an acetylene affords a silver mirror, and the cations [Mo(CO) $\left(\eta^{2}-\mathrm{RC}_{2} \mathrm{R}^{1}\right)_{2}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right.$ or $\left.\left.\eta^{5}-\mathrm{C}_{9} \mathrm{H}_{7}\right)\right]^{+} \mathrm{BF}_{4}^{-}\left(R=R^{1}=H ; R=R^{1}=\mathrm{Me} ; R=\mathrm{Bu}-t, R^{1}=H ; R=\mathrm{Me}, R^{1}\right.$ $=H ; R=P h, R^{1}=H ; R=R^{1}=P h$ ), which on treatment (room temperature, $\mathrm{CH}_{2} \mathrm{Cl}_{2}$ ) with $\mathrm{P}(\mathrm{OMe})_{3}$ (excess) gives the highly colored cations 1 and 2 in good yield. The HNMR spectrum $\left(20^{\circ} \mathrm{C}, \mathrm{CD}_{3} \mathrm{NO}_{2}\right)$ of, for example, $\mathrm{Mo}\left(\eta^{2}-\right.$ $\left.\left.\mathrm{MeC}_{2} \mathrm{Me}\right)\left[\mathrm{P}(\mathrm{OMe})_{3}\right]_{2}\left(\eta^{5}-\mathrm{C}_{9} \mathrm{H}_{7}\right)\right]^{+} \mathrm{BF}_{4}{ }^{-}$shows a sharp triplet $\left({ }^{4} \mathrm{~J}_{\mathrm{HP}}=1.0 \mathrm{~Hz}\right.$.) at $\tau 7.7$ ( MeCOC ) collapsing reversibly (coalescence temperature, in $\left(\mathrm{CD}_{3}\right)_{2} \mathrm{CO}$ at $\left.-40^{\circ} \mathrm{C}\right)$ to two resonances ( $\tau 7.3(3 \mathrm{H}, \mathrm{brs}), 8.4(3 \mathrm{H}, \mathrm{br} \mathrm{s})$ ) on cooling, which suggests that acetylene propellor rotation occurs; the observation of ${ }^{31} \mathrm{P}$ coupling excludes acetylene dissociation. The coalescence temperature of the corresponding $\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}$ substituted cation is even lower $\left(>-90^{\circ} \mathrm{C}\right.$ ), which is interesting in view of the report ${ }^{2}$ that the ${ }^{1} \mathrm{H}$ NMR spectrum $\left(\left(\mathrm{CD}_{3}\right)_{2} \mathrm{CO}\right)$ of $\left[\mathrm{Mo}\left(\eta^{2}-\mathrm{MeC}_{2} \mathrm{Me}\right)(\text { diphos })\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)\right]^{+} \mathrm{PF}_{6}{ }^{-}$collapses to a single $\mathrm{MeC}=\mathrm{C}$ signal only at $110^{\circ} \mathrm{C}$.
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## The Curious Structure of the Lithiocarbon ${ }^{1} \mathbf{C}_{3} \mathrm{Li}_{4}$

Sir:
$\mathrm{C}_{3} \mathrm{Li}_{4}$, prepared by lithiation of propyne by $n$-butyllithium, ${ }^{2 d, 3}$ is a readily available lithiocarbon. ${ }^{1}$ Only the most general information is available concerning the structure of this species. $\mathrm{C}_{3} \mathrm{Li}_{4}$ reacts to give either allene- or propyne-type products depending on the reagents and the conditions. ${ }^{2 \mathrm{~d}}$ This suggests an open rather than a cyclic arrangement of the carbon atoms. The IR absorption band of $\mathrm{C}_{3} \mathrm{Li}_{4}$ in the $1700-\mathrm{cm}^{-1}$ region has been interpreted in terms of formulation I or II. ${ }^{2 d, 3}$


We have used ab initio molecular orbital calculations ${ }^{4}$ to investigate possible $\mathrm{C}_{3} \mathrm{Li}_{4}$ structures. Complete geometry optimizations within each assumed symmetry were carried out using the minimal STO-3G basis set. ${ }^{4 \mathrm{~b}}$ Single point calculations using these optimized geometries followed, employing the split valence $4-31 \mathrm{G}$ (for carbon) ${ }^{4 \mathrm{c}}$ and 5-21G (for lithium ${ }^{4 \mathrm{~d}}$ bases. Our earlier work has shown that polylithio derivatives adopt highly unconventional geometries; ${ }^{5}$ the same trend was expected in the present search. Nevertheless, we began with the lithiated a nalogue of allene, III. Although III proved to be a local minimum when optimized within the constraint of $D_{2 d}$ symmetry, the energy compared unfavorably with that of other structures. The planar $D_{2 h}$ allene (IV) (with four $\pi$ electrons), which had nearly the same bond lengths and angles, was more stable by $0.5 \mathrm{kcal} / \mathrm{mol}$ at the $4-31 \mathrm{G} / 5-21 \mathrm{G}$ level (Table I). The theoretical estimates for the rotational barrier in allene are between 27 and $92 \mathrm{kcal} / \mathrm{mol}^{1}{ }^{6}$ experimental values of $\sim 46 \mathrm{kcal} / \mathrm{mol}$ have been reported for $1,3-$ dialkylallenes. ${ }^{7}$ Lithiation is thus able not only to reduce the $D_{2 d}-D_{2 h}$ energy differences dramatically, but also to reverse the normal order of stability. Similarly, 1,1-dilithioethylene has been found to prefer the perpendicular, rather than the normal planar conformation. ${ }^{5 \mathrm{a}}$

Encouraged by the tendency of lithium to bridge, ${ }^{5 \mathrm{c}}$ we examined the tetrabridged $D_{2 d}$-constrained structure, V , which proved to be more stable than III by $19.5 \mathrm{kcal} / \mathrm{mol}$ at STO-3G but $13.3 \mathrm{kcal} / \mathrm{mol}$ less stable at $4-31 \mathrm{G} / 5-21 \mathrm{G}$. Since both III and V are $D_{2 d}$ it is evident that a barrier (within this symmetry) exists between them. On the other hand, the conventional acetylenic structure, VI, is not a local minimum, but collapses to a triply bridged form, VII, upon optimization under $C_{3 v}$ symmetry. VII is more stable than III-V at both basis set levels. The effect of the three bridging lithiums is seen in the

Table I. Structure and Relative Energies of Various $\mathrm{C}_{3} \mathrm{Li}_{4}$ Isomers

| Molecule | Point group | Parameter | Values ${ }^{\text {a }}$ | Relative energy, kcal mol |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  | $\begin{aligned} & \text { STO-3G } \\ & \text { opt. } \end{aligned}$ | $\begin{gathered} 4-31 \mathrm{G} / 5- \\ 21 \mathrm{G} \end{gathered}$ |
| III | $D_{2 d}$ | $r\left(\mathrm{C}_{1}-\mathrm{C}_{2}\right)$ | 1.323 | 35.6 | 19.0 |
|  |  | $r\left(\mathrm{Li}_{4}-\mathrm{C}_{1}\right)$ | 1.791 |  |  |
|  |  | $\angle \mathrm{Li}_{4} \mathrm{C}_{1} \mathrm{C}_{2}$ | 122.3 |  |  |
| IV | $D_{2 h}$ | $r\left(\mathrm{C}_{1}-\mathrm{C}_{2}\right)$ | 1.322 | 34.2 | 18.5 |
|  |  | $r\left(\mathrm{Li}_{4}-\mathrm{C}_{1}\right)$ | 1.807 |  |  |
|  |  | $\angle \mathrm{Li}_{4} \mathrm{C}_{1} \mathrm{C}_{2}$ | 122.6 |  |  |
| V | $D_{2 d}$ | $r\left(\mathrm{C}_{1}-\mathrm{C}_{2}\right)$ | 1.371 | 16.1 | 32.3 |
|  |  | $r\left(\mathrm{C}_{1}-\mathrm{Li}_{4}\right)$ | 1.902 |  |  |
|  |  | $r\left(\mathrm{C}_{2}-\mathrm{Li}_{4}\right)$ | 1.928 |  |  |
| VII | $C_{3 c}$ | $r\left(\mathrm{C}_{1}-\mathrm{C}_{2}\right)$ | 1.409 | 14.3 | 12.4 |
|  |  | $r\left(\mathrm{C}_{2}-\mathrm{C}_{3}\right)$ | 1.279 |  |  |
|  |  | $r\left(\mathrm{C}_{3}-\mathrm{Li}_{4}\right)$ | 1.774 |  |  |
|  |  | $r\left(\mathrm{C}_{1}-\mathrm{Li}_{5}\right)$ | 1.913 |  |  |
|  |  | $\angle \mathrm{Li}_{5} \mathrm{C}_{1} \mathrm{C}_{2}$ | 69.6 |  |  |
| VIII | $D_{2 h}$ | $r\left(\mathrm{C}_{1}-\mathrm{C}_{2}\right)$ | 1.316 | 22.6 |  |
|  |  | $r\left(\mathrm{Li}_{4}-\mathrm{C}_{1}\right)$ | 1.741 |  |  |
|  |  | $r\left(\mathrm{C}_{2}-\mathrm{Li}_{5}\right)$ | 1.830 |  |  |
| $1 \mathrm{X}^{d}$ | $C_{2 c}$ | $r\left(\mathrm{C}_{1}-\mathrm{C}_{2}\right)$ | 1.326 | $0.0^{\text {b }}$ | $0.0{ }^{\text {c }}$ |
|  |  | $r\left(\mathrm{C}_{1}-\mathrm{Li}_{4}\right)$ | 1.794 |  |  |
|  |  | $r\left(\mathrm{C}_{2}-\mathrm{Li}_{5}\right)$ | 1.869 |  |  |
|  |  | $\angle \mathrm{Li}_{4} \mathrm{C}_{1} \mathrm{C}_{2}$ | 149.3 |  |  |
|  |  | $\angle C_{1} C_{2} C_{3}$ | 155.7 |  |  |
|  |  | $\angle \mathrm{Li}_{5} \mathrm{C}_{2} \mathrm{Li}_{6}$ | 93.2 |  |  |

${ }^{a}$ STO-3G optimized geometries, bond lengths in Ångströms and bond angles in degrees. ${ }^{\text {b }}$ Corresponding total energy. -141.44166 au. corresponding total energy, -143.08011 au. ${ }^{d}$ See IX for specification of chain structure.
$\mathrm{C}-\mathrm{C}$ bond distances of VII; the $\mathrm{C}-\mathrm{C}$ triple bond is lengthened to $1.279 \AA$ while the single bond is shortened to $1.409 \AA$ (in propyne $C_{1}-C_{2}=1.170 \AA$ ). ${ }^{8}$ Lithium substitution also lengthens the multiple bonds in III, IV, and V, where the $\mathrm{C}-\mathrm{C}$ lengths are considerably longer than the calculated RHF/ STO-3G value for allene ( $1.288 \AA$ ). ${ }^{8}$

Although having no precedent in normal hydrocarbon structures, VIII ( $D_{2 h}$ ) with a planar tetracoordinate carbon ${ }^{56}$ is found to be remarkably low in energy. Distortion of VIII to $C_{2 t}$ symmetry by decreasing the $\mathrm{Li}_{5} \mathrm{C}_{2} \mathrm{Li}_{6}$ angle destabilizes the system if the atoms around $\mathrm{C}_{2}$ are moved in opposite directions toward a tetrahedral arrangement. (Such distortion ultimately would lead to a tetralithiocyclopropene structure. We also examined a face lithiated $C_{3 i}$ tetralithiocyclopropene which was similarly quite uncompetitive energetically.) On the other hand movement of the atoms around $\mathrm{C}_{2}$ in the same direction (also $C_{2 c}$ ) so that $\mathrm{C}_{1}, \mathrm{C}_{3}, \mathrm{Li}_{5}$, and $\mathrm{Li}_{6}$ lie in the same hemisphere (IX) gives the lowest energy $\mathrm{C}_{3} \mathrm{Li}_{4}$ structure that we have been able to find (both at STO-3G and at 4-31G/521 G -see Table I). A further reduction of symmetry to $C_{2}$ by moving the central lithiums toward opposite carbons resulted in destabilization. IX possesses a zig-zag Li-C-C-C-Li chain, but one not deviating far from linearity. However, the two additional lithium atoms, $\mathrm{Li}_{5}$ and $\mathrm{Li}_{6}$. lie in mutually perpendicular positions $\left(\angle \mathrm{Li}_{5} \mathrm{C}_{2} \mathrm{Li}_{6}, 93.2^{\circ}\right)$. The stability of this arrangement can be understood by considering it in a formal way as a $\mathrm{C}_{3}{ }^{4-}$ ion stabilized by four surrounding $\mathrm{Li}^{+}$cations. The linear $\mathrm{C}_{3}{ }^{4-}$ fragment is isoelectronic with carbon dioxide and will have allylic type stabilization in both mutually perpendicular $\pi$ systems. The terminal Li ions, $\mathrm{Li}_{4}$ and $\mathrm{Li}_{7}$, form partially covalent $\sigma$ bonds with $\mathrm{C}_{1}$ and $\mathrm{C}_{3}$ and also act as $\pi$ acceptors. The central Li ions, $\mathrm{Li}_{5}$ and $\mathrm{Li}_{6}$, interact favorably with highest occupied nonbonding $\pi$ orbitals of the linear $\mathrm{C}_{3}$ chain primarily as $\pi$ acceptors (X). ${ }^{9}$ Clearly, the nearly perpendicular angle for $\mathrm{Li}_{5} \mathrm{C}_{2} \mathrm{Li}_{6}$ allows such interactions for both allylic nonbonding $\pi$ orbitals and thus gives a lower energy than VIII in which only one such orbital can be stabilized. The

Mulliken population analysis supports this interpretation. All lithium atoms have substantial valence orbital populations, but, for $\mathrm{Li}_{5}$ and $\mathrm{Li}_{6}$, these are primarily associated with the p functions with axes parallel to the $\mathrm{C}_{1} \mathrm{C}_{3}$ axis. The ion-pair
 have stressed in earlier papers ${ }^{5}$ and we reiterate here that it is more satisfactory to emphasize the multicenter, covalent bonding of such polylithium compounds rather than their presumed "ion-pair" character.

The electropositive nature of lithium, and the availability of low lying p orbitals, results in lithiocarbon structures which are quite different from those of the corresponding hydrocarbons. The "bare" carbon ${ }^{10}\left(\mathrm{C}_{2}\right)$ of IX is remarkable; all four valencies extend in one direction in roughly $C_{4 v}$ fashion. Bare carbons are also possessed by $V\left(C_{1}\right.$ and $C_{3}$ are tricoordinate, each with local $C_{2 t}$ symmetry) and by VII ( $\mathrm{C}_{1}$, tetracoordinate, has $C_{3 t}$ symmetry), but these structures, less stable than IX, may not be capable of existence.

In theoretical structural searches of this type where conventional experience is no guide, it is difficult to establish with certainty that the absolute minimum has been found. The large number of geometrical possibilities, and the time required for these calculations ( $\sim 45 \mathrm{~min}$ for each STO-3G C ${ }_{3} \mathrm{Li}_{4}$ point on the Telefunken TR 440), precludes a global search. We did not, for example, examine any structures with $C_{s}$ or with $C_{1}$ point groups. Results at the Hartree-Fock level with the minimal STO-3G or with the split valence $4-3 \mathrm{IG} / 5-21 \mathrm{G}$ bases cannot be considered to be definitive, although (with the exception of V) the same relative energy orderings are found (Table I). Nevertheless, our previous experience with calculations on polylithio derivatives at higher levels including correlation effects indicates that the procedure used here may give reasonable energy orderings. ${ }^{5}$ In particular, the relative energies of linear and bridged dilithioacetylene were not greatly changed by allowance of electron correlation. ${ }^{5 \mathrm{c}}$ As always, such calculations refer to isolated species in a vibrationless state. However, $\mathrm{C}_{3} \mathrm{Li}_{4}$ can be prepared in a hydrocarbon solvent, and may not be highly associated. ${ }^{2 \mathrm{~d}, 3}$ The ${ }^{7} \mathrm{Li}$ and ${ }^{13} \mathrm{C}$ NMR spectra should be of interest, to say nothing of more refined experimental structural information. We will report the unusual structures of other lithiocarbons in due course. ${ }^{12}$

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## A Four-Center, Concerted, Bimolecular Reaction: $\mathrm{ICl}{ }^{*}\left(\mathbf{A}^{3} \Pi_{1}\right)+\mathrm{H}_{\mathbf{2}} \rightarrow \mathbf{H C l}+\mathbf{H I}$

Sir:
The rate of the laser-induced chemical reaction

$$
\begin{equation*}
\mathrm{ICl} *\left(\mathrm{~A}^{3} \Pi_{1}\right)+\mathrm{H}_{2} \rightarrow \mathrm{HCl}+\mathrm{HI} \tag{1}
\end{equation*}
$$

has been measured as a function of ICl * vibrational energy. There is no dark reaction, but the reaction is promoted by water vapor, stopcock grease, metal surfaces, and stray light. ${ }^{1}$ The kinetics appears similar to $\mathrm{I}_{2}{ }^{*}+\mathrm{H}_{2},{ }^{2-4}$ and other reactions have been reported which involve electronic excitation of $\mathrm{I}_{2}{ }^{5}$ The problem of four-center reactions has received considerable attention. ${ }^{6}$

ICl is prepared under vacuum from $\mathrm{I}_{2}$ and $\mathrm{Cl}_{2}{ }^{7,8}$ and is never exposed to air. Mixtures of $\mathrm{ICl}\left(\sim 5\right.$ or 9 Torr) and $\mathrm{H}_{2}$ (10-600 Torr) are photolyzed for about 90 min in a $1-\mathrm{m}$-long, $5-\mathrm{cm}$ diameter quartz cell by a Chromatix CMX-4 flashlamppumped dye laser ( $\Delta \nu=0.3 \mathrm{~cm}^{-1}$ ) run at 30 pps . A calibrated thermopile monitors the laser power, which is 1,5 to 6 mJ per pulse in a $1-\mathrm{cm}^{2}$ beam. Approximately $60 \%$ absorption of the laser light is measured at $\mathrm{I}^{35} \mathrm{Cl} v^{\prime \prime}=0$ bandheads.
After photolysis all of the HCl , together with some residual $\mathrm{Cl}_{2},{ }^{7}$ is distilled into a bulb with $\mathrm{I}_{2}$ crystals. The $\mathrm{Cl}_{2}$ reacts completely with the $\mathrm{I}_{2}$, and the HCl is then redistilled. Its pressure is measured with a capacitance manometer, and mass spectra confirm that the product is HCl . HI and ICl react to form HCl and $\mathrm{I}_{2}$; thus two molecules of HCl are formed in each reaction.
$\mathrm{ICl} *$ is removed by reaction 1 and by quenching:


Figure 1. Solid points are data without $\mathrm{Ar}, P_{1 \mathrm{C} 1}=8.9$ Torr; the solid line is a fit to these points. The dashed line is the predicted yield with $P_{\text {Ar }} / P_{\mathrm{H}_{2}}$ $=4.75$; open circles are corresponding data. Here, $P_{1 \mathrm{Cl}}=4.9$ Torr except for lowest $P_{\mathrm{H}_{2}} / P_{1 \mathrm{Cl}}$ point, where $P_{1 \mathrm{Cl}}=8.9$ Torr ( $\lambda 616.8 \mathrm{~nm}$, exciting $1200 \mathrm{~cm}^{-1}$ below dissociation).


Figure 2. Quantum yield vs. excitation energy. $P_{1 \mathrm{CI}}=8.9$ Torr. The solid line is a fit to the data; the dashed line is a rough estimate of the effect of collisional dissociation. The point at $18800 \mathrm{~cm}^{-1}$ was taken using a doubled YAG laser with $60-\mathrm{mJ}, 20-\mathrm{ns}$ pulses.

$$
\begin{align*}
\mathrm{ICl}^{*}+\mathrm{ICl} & \rightarrow \mathrm{ICl}+\mathrm{ICl}  \tag{2}\\
\mathrm{ICl}^{*}+\mathrm{H}_{2} & \rightarrow \mathrm{ICl}+\mathrm{H}_{2} \tag{3}
\end{align*}
$$

Fluorescence is much too slow to complete with reactions 1-3. The quantum yield for reaction is

$$
\begin{equation*}
\Phi=\frac{(\mathrm{HCl} \text { produced }) / 2}{\text { photons absorbed }}=\frac{k_{1}\left[\mathrm{H}_{2}\right]}{k_{2}[\mathrm{ICl}]+\left(k_{1}+k_{3}\right)\left[\mathrm{H}_{2}\right]} \tag{4}
\end{equation*}
$$

A Stern-Volmer plot of ICl* fluorescence lifetime vs. $\mathrm{H}_{2}$ pressure gives $\left(k_{1}+k_{3}\right)=1.5 \times 10^{-11} \mathrm{~cm}^{3}$ molecule ${ }^{-1} \mathrm{~s}^{-1}$; similarly, Steinfeld ${ }^{9}$ has measured $k_{2}=2.3 \times 10^{-10} \mathrm{~cm}^{3}$ molecule ${ }^{-1} \mathrm{~s}^{-1}$; both are nearly independent of wavelength. $k_{1}$ is determined from the quantum yield, Figure 1. The solid curve is a least-square fit to the data using eq 4 . Note that a two-parameter curve is fit with a single free parameter. The result, $k_{1}=(9.0 \pm 0,8) \times 10^{-14} \mathrm{~cm}^{3}$ molecule ${ }^{-1} \mathrm{~s}^{-1}(90 \%$ confidence limits), corresponds to 9000 gas kinetic collisions. Assuming that the preexponential factor does not exceed the gas kinetic collision frequency, the activation energy is less than $5.5 \mathrm{kcal} / \mathrm{mol}$. The reaction rate is quite large; the quantum yield is low because the quenching rate is even larger.

A single vibrational level of ICl * is excited initially, but vibrationally inelastic collisions spread the population distribution so that reaction may occur from other levels. The extent of this spread, which is limited by quenching, can be estimated

